

# A new model for fully anisotropic diffusion of CR protons and electrons in the heliosphere

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### Anisotropic Diffusion of Energetic Particles in Heliospheric and Galactic Magnetic Fields

Two faces of the same coin:

## **Cosmic Ray modulation in the Heliosphere**



### **Cosmic Ray transport in the Galaxy**

$$\frac{\partial f}{\partial t} = \nabla \cdot (\hat{\kappa} \,\nabla f) - \vec{v} \cdot \nabla f + p \left(\frac{1}{3} (\nabla \cdot \vec{v}) + a_{\pi}\right) \frac{\partial f}{\partial p} + 3a_{\pi}f + \frac{q}{p^2}$$



### Anisotropic Diffusion of Energetic Particles in Heliospheric and Galactic Magnetic Fields



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### **Galactic Protons**



$$j = 12.14 \cdot \beta (E_k + 0.5E_0)^{-2.6}$$
  
Bow Shock  
Heliospheric Shock  
Local  
Interstellar



[Reinecke et al., 1993, JGR]

### **Jovian Electrons**

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### On the propagation of Jovian electrons in the heliosphere: transport modelling in 4-D phase space

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$$j_{gal}(R(p)) = \delta \exp \left( \alpha (\ln R)^2 + \beta \ln R + \gamma \right) (5)$$
  
$$j_{Jov}(E(p)) = j(E_r) \frac{E_r^{1.5} (h + E_r)^{3.5}}{E^{1.5} (h + E)^{3.5}}$$
(6)





### Two Questions, Two Complementary Numerical Tools

Heliophysical Questions:

- How are the cosmic ray protons and electrons modulated on their way from the heliospheric boundary to earth?
- Can we determine transport parameters from the propagation of Jovian electrons?
- Grid-Based Numerics: DuFort-Frankel, VLUGR3 ...

Solving the CR-Transport Equation directly on a numerical grid via finite differences or finite volume methods.

• Stochastic Differential Equations (SDE)

Solve a SDE equivalent to the Transport Equation by tracing pseudo-particles

# Testing the SDE Modulation with an Analytic Parker Propagator

### The Parker propagator for spherical solar modulation

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$$\begin{split} j(r,p) &= p^2 f(r,p) \\ &= \frac{3}{b_p} \int dr_0 \int dp_0 \frac{S(r_0,p_0)}{V(r_0)} \frac{p_0 y_0}{f_p} \left(\frac{r_0}{r}\right)^{\frac{1+\beta_p}{2}} \\ &\times \left(\frac{p_0}{p}\right)^{\frac{3\beta_p - 4\alpha_p - 5}{2(2+\alpha_p)}} \exp\left(-\frac{y_0(1+h_p^2)}{f_p}\right) \\ &\times I_{\frac{1+\beta_p}{1+\alpha_p - \beta_p}} \left(\frac{2y_0 h_p}{f_p}\right) \end{split}$$



**Fig. 1.** The modulated spectra of ACRs and GCRs in the heliosphere. The solar wind termination shock marking the position of the sources is located at  $r_{\rm sh} = 100$  AU. The solid lines are the combined spectra, the dotted and dashed lines indicate the individual contributions from ACRs and GCRs, respectively.

### **Results** Testing the SDE Modulation with an Analytic Parker Propagator



$$\begin{aligned} \frac{\partial f}{\partial t} &= \nabla \cdot \left( \hat{\kappa} \, \nabla f \right) \\ \text{local Tensor} \\ \hat{\kappa}_L &= \begin{pmatrix} \kappa_{\perp 1} & 0 & 0 \\ 0 & \kappa_{\perp 2} & 0 \\ 0 & 0 & \kappa_{\parallel} \end{pmatrix} \end{aligned}$$





κ̂<sub>G</sub>

global RF

global RF







### **Anisotropic Diffusion**

The "Classical" Approach: Euler-Angle Transformation

$$\kappa_{rr} = \kappa_{\perp,2} \sin^2 \zeta + \cos^2 \zeta \left( \kappa_{||} \cos^2 \Psi + \kappa_{\perp,3} \sin^2 \Psi \right),$$
  

$$\kappa_{r\theta} = -\kappa_A \sin \Psi + \sin \zeta \cos \zeta$$
  

$$\times \left( \kappa_{||} \cos^2 \Psi + \kappa_{\perp,3} \sin^2 \Psi - \kappa_{\perp,2} \right),$$
  

$$\kappa_{r\phi} = -\kappa_A \cos \Psi \sin \zeta - \left( \kappa_{||} - \kappa_{\perp,3} \right) \sin \Psi \cos \Psi \cos \zeta,$$
  

$$\kappa_{\theta r} = \kappa_A \sin \Psi + \sin \zeta \cos \zeta$$
  

$$\times \left( \kappa_{||} \cos^2 \Psi + \kappa_{\perp,3} \sin^2 \Psi - \kappa_{\perp,2} \right),$$
  

$$\kappa_{\theta \theta} = \kappa_{\perp,2} \cos^2 \zeta + \sin^2 \zeta \left( \kappa_{||} \cos^2 \Psi + \kappa_{\perp,3} \sin^2 \Psi \right),$$
  

$$\kappa_{\theta \phi} = \kappa_A \cos \Psi \cos \zeta - \left( \kappa_{||} - \kappa_{\perp,3} \right) \sin \Psi \cos \Psi \sin \zeta,$$
  

$$\kappa_{\phi r} = \kappa_A \cos \Psi \sin \zeta - \left( \kappa_{||} - \kappa_{\perp,3} \right) \sin \Psi \cos \Psi \sin \zeta,$$
  

$$\kappa_{\phi \theta} = -\kappa_A \cos \Psi \cos \zeta - \left( \kappa_{||} - \kappa_{\perp,3} \right) \sin \Psi \cos \Psi \sin \zeta,$$
  

$$\kappa_{\phi \phi} = \kappa_{||} \sin^2 \Psi + \kappa_{\perp,3} \cos^2 \Psi,$$
  
(17)

where  $\tan \Psi = -B_{\phi}/(B_r^2 + B_{\theta}^2)^{1/2}$  and  $\tan \zeta = B_{\theta}/B_r$ . Note that

### Anisotropic Diffusion The new Approach: Local Frenet-Trihedron

Tangential Vector 
$$\vec{t} = \frac{\vec{B}}{B}$$

$$(\vec{t}\cdot\nabla)\vec{t} = k\vec{n} \; ; \; (\vec{t}\cdot\nabla)\vec{n} = -k\vec{t} + \tau\vec{b} \; ; \; (\vec{t}\cdot\nabla)\vec{b} = -\tau\vec{n}$$







The symmetric global diffusion tensor:

$$\hat{\kappa}_{11} = \kappa_{\perp 1}n_1^2 + \kappa_{\perp 2}b_1^2 + \kappa_{\parallel}t_1^2 
\hat{\kappa}_{12} = \kappa_{\perp 1}n_1n_2 + \kappa_{\perp 2}b_1b_2 + \kappa_{\parallel}t_1t_2 
\hat{\kappa}_{13} = \kappa_{\perp 1}n_1n_3 + \kappa_{\perp 2}b_1b_3 + \kappa_{\parallel}t_1t_3 
\hat{\kappa}_{22} = \kappa_{\perp 1}n_2^2 + \kappa_{\perp 2}b_2^2 + \kappa_{\parallel}t_2^2 
\hat{\kappa}_{23} = \kappa_{\perp 1}n_2n_3 + \kappa_{\perp 2}b_2b_3 + \kappa_{\parallel}t_2t_3 
\hat{\kappa}_{33} = \kappa_{\perp 1}n_3^2 + \kappa_{\perp 2}b_3^2 + \kappa_{\parallel}t_3^2$$

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# General magnetic field & local diffusion coefficients

The symmetric global diffusion tensor:

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In any case: The simple structure of the diffusion tensor in the local frame may lead to complicated tensor elements in the global frame, depending on the magnetic field.



### RUB

### **Anisotropic Diffusion** The Parker-Field Tensor Elements (isotropic perp.)

$$\begin{cases} \kappa_{\parallel} = 1, \ \kappa_{\perp 1} = 0.2, \ \kappa_{\perp 2} = 0.2 \\ r = 10 \text{AU}, \ \varphi = \pi \end{cases}$$



Colatitude

### RUB

### **Anisotropic Diffusion** The Parker-Field Tensor Elements (anisotropic perp.) $\kappa_{\parallel} = 1, \, \kappa_{\perp 1} = 0.05, \, \kappa_{\perp 2} = 0.2$ $r = 10 \text{AU}, \varphi = \pi$ 1.2 (old) k11 k1 (old) k13 k13 (old) k12 k12 1 0.8 Tensor elements 0.6 0.4 0.2 0 -0.2 -0.4 1.2 (old) k22 k22 (old) k33 (old) k23 k23 1 0.8 Tensor elements 0.6 0.4 0.2 0 -0.2 -0.4 20 40 60 80 100120140160 20 40 60 80 100120140160 20 40 60 80 100120140160

Colatitude

Colatitude

### **Results** Modulation Spectra Isotropic vs Anisotropic



### **Results** Modulation Spectra Isotropic vs Anisotropic (Normalized to LIS)



### **Results** Effect of Different Tensor Formulations

Isotropic perpendicular diffusion 1 AU Earth spectra Anisotropic perpendicular diffusion 1 AU Earth spectra



80

60

40

20

0

-20

-40

-60

-80

-100

Z-Axis

### **Results** Jovian Electron Distribution (8 MeV)



With Standard Modification:

-50

$$F( heta) = -C^+ + C^- anh \left[ ( ilde{ heta} - 90^\circ - heta_F) / \Delta heta 
ight]$$

Isotropic Perpendicular Diffusion



With Standard Modification:

$$F( heta) = -C^+ + C^- anh \left[ ( ilde{ heta} - 90^\circ - heta_F) / \Delta heta 
ight]$$

Isotropic Perpendicular Diffusion





Isotropic Perpendicular Diffusion

Anisotropic Perpendicular Diffusion:  $\kappa_{\perp 2} = 0.5 \kappa_{\perp 1}$ 



Anisotropic Perpendicular Diffusion:  $\kappa_{\perp 2} = 0.1 \kappa_{\perp 1}$ 

Isotropic Perpendicular Diffusion

### Conclusions

- Anisotropic Diffusion is important for Cosmic Ray transport on different scales and its proper treatment is needed to determine their actual distribution in the Galaxy and the modulation effects in the Heliosphere.
- Complementary numerical tools exist and are under development to investigate the properties of solutions to the Parker transport equation for various coordinate systems, setups and boundary conditions.
- Our new approach to the formulation of the diffusion tensor results in differences for the tensor elements in the fully anisotropic case and has a possible impact on the resulting modulation spectra, especially for high latitudes.
- The LIS modulation of galactic protons and the distribution of Jovian electrons shows to be sensitive to the actual structure of the diffusion tensor.