MHD modeling of the inner heliosphere and its transients

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Talk outline:

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[Motivation for studying \(I\)CMEs](#page-5-0)

Coronal mass ejection (CME): large blob of solar plasma ($m \approx 10^{13}$ kg, $v_0 \approx 20...3000$ km/s) ejected spacewards

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Why care about CMEs?

- **1** Major manifestation of solar activity
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	- particle acceleration at shocks
	- global flux removal, ...
- - safety concerns for astronautics,
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- 2 CMEs relate to many other fields of solar physics
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	- particle acceleration at shocks
	- global flux removal, ...
- 3 Commercial application: **"space weather"**
	- safety concerns for astronautics,
	- satellite communication failures, etc.
	- \Rightarrow urgent need to predict outbreak and IP evolution!

- **1** The CME phenomenon spans vast temporal and spatial scales. ⇒ Need to specialize on selected aspects/phases.
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What makes CME modeling a demanding task?

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Focus of MHD-based CME models can be on

- initiation/eruption [not this talk]
- propagation (expansion, trajectory, acceleration/travel time)
- • interaction (with the solar wind/CIRs, other CMEs, and/or planetary magnetospheres)

[Self-consistent \(MHD\) modeling](#page-6-0) [Numerics and physical realization](#page-13-0)

(Simplifying) analytical CME models are few in number. Space weather prediction relies on large-scale numerical MHD.

CSEM [Tóth 2005] CISM [Odstrcil 2008]

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Major (technical) challenge: High resolution requirements due to

- **1** need to track features $\ll R_{\odot}$ across $>$ 214 $R_{\odot} = 1$ AU
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- **1** need to track features $\ll R_{\odot}$ across $>$ 214 $R_{\odot} = 1$ AU
- **2** Lack of symmetry

solar min: B_0 is 2D, but CME expansion $\frac{p}{q}$ dipolar axis solar max: **B**_o is 3D itself

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Solution #1: Ignore φ dependence anyway.

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Solution #2: Performance tuning

- specially tailored grids, esp. spherical with radially varying ∆*r* = ∆*r*(*r*)
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- multi-scale models [e.g. Riley et al. 2006] (NB: $\|u\| > v_A$ after a few R_{\odot})

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Existing models can be classified by...

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- **4** Realization of boundary conditions at $r = R_0$ (analytic
	- uniform [e.g. Vandas et al. 1998, 2002],
	- structured [Odstrcil & Pizzo 1999; Manchester et al. 2004],
	- **realistic** [Hayashi et al. 2006; Shen et al. 2007]

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- **Realization of boundary conditions at** $r = R_0$ **(analytic** or observationally derived, e.g. from magnetograms) and the background solar wind
	- uniform [e.g. Vandas et al. 1998, 2002],
	- structured [Odstrcil & Pizzo 1999; Manchester et al. 2004],
	- realistic [Hayashi et al. 2006; Shen et al. 2007]

- **5** Included physics, e.g. treatment of the energy budget:
- isothermal $\gamma = 1$ or adiabatic $\gamma = \gamma_0 \leq 5/3$ ($\rho \sim \rho^{\gamma}$)
- $\bullet \quad \gamma = \gamma(r)$ [e.g. Fahr et al. '76, Lugaz et al. '07] (\blacksquare not good for shocks)
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	- 2 $S = u_1 \cdot \nabla (p_1 \rho_1^{-5/3})$ $\binom{-5/3}{1}$ with $(\cdot)_1$ from $\gamma=1.05$ run [Pomoell, Vainio, Kissmann 2011]
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	- **3** Consistent Alfvénic wave heating.

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SW heating by Alfvén waves

Concept:

Waves are excited near R_o, travel along **B**, get shifted up in *f*, and dissipate at *f*h. Variables: either scalar fields ε_{+} or full spectrum $P(f, r, t)$.

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 $\partial_t P + \nabla \cdot [(\mathbf{u} \pm \mathbf{v}_A) P] + (P/2) \nabla \cdot \mathbf{u} = -\partial_t F$ gives

- wave pressure $p_w(\mathbf{r}) = (1/2) \int_{t_0}^{t_h(\mathbf{r})} P(f, \mathbf{r}) \, \mathrm{d}f \triangleq (\varepsilon_+ + \varepsilon_-)/2$
- heating term $Q_w = F(f_h, r) P(f_h(r), r)$ [**u** $\pm \mathbf{v}_A$] · $\nabla f_h(r)$

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[Some principal findings](#page-35-0) [Connecting to observations](#page-36-0) [Summary/ Conclusions](#page-38-0)

Some findings from principal models

- **1** Some indication of approximately self-similar evolution [e.g. Kleimann et al. '09]
- - background SW (higher speeds in fast,
	- the initial polarity w.r.t. B_{sw} , influencing

[Some principal findings](#page-34-0) [Connecting to observations](#page-36-0) [Summary/ Conclusions](#page-38-0)

Some findings from principal models

- **1** Some indication of approximately self-similar evolution [e.g. Kleimann et al. '09]
- **2** CME development strongly depends on
	- background SW (higher speeds in fast, dilute winds) [Jacobs et al. 2005] and
	- the initial polarity w.r.t. \mathbf{B}_{sw} , influencing the CME's speed, shape, and deflection.

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Comparison to (satellite) observations

Lugaz et al. [2007] Chané et al. [2008]

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Comparison to (satellite) observations

Lugaz et al. [2007] Chané et al. [2008]

SWMF (3D, MDI init) VAC (2.5D, analyt. init)

NB: Models are quite sensitive to chosen parameters [e.g. Schrijver et al. 2008], but published results often consider only limited parameter ranges.

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Conclusions

- CMEs are a very diverse class of heliospheric transients.
- 3D MHD simulations are indispensable to model a CME's life cycle, with the long-term goal of reliable forecasts.
- Modeling results/predictions crucially depend on initial parameters and physical effects included.
- Models benefit from high-quality S/C data input to
	- **1** constrain IC/BCs and
	- 2 allow for a posteriori verification of results.
- • Simple models can be useful, provided their limitations are taken into account. \rightarrow Importance of comparative studies!